Pulsar Wind Nebulae as High Energy γ -Ray Emitters

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Abstract. There is growing evidence that some variable high-energy γ -ray sources are associated with pulsar wind nebulae. We review the current status of X-ray and radio studies of the sources most likely to be significant emitters above 100 MeV. The assumption of the γ -radiation arising from synchrotron processes puts a lower limit on the magnetic field and suggests there may be significant doppler boosting causing the asymmetries seen in the morphology of these nebulae. However, these sources are also in close proximity to molecular clouds, suggesting that variable γ -ray emission may arise from a ram pressure confined pulsar wind nebula in a dense medium.

1. Introduction

For over 30 years, the nature of the majority of high-energy $(E>100~{\rm MeV})$ γ -ray sources in the Galactic plane has been a mystery. The only firmly identified source class has been spin-powered pulsars, whose phase-averaged emission appears constant. However, the EGRET data clearly show that there is at least one class of variable emitters in the Galaxy (eg. Tompkins 1999). Several suggestions as to the nature of these objects have been put forward but no strong observational evidence has previously supported a particular source class.

However, the variability timescale of a few months to a few years for sources with emission above 100 MeV gives us some clues to their nature. First, a maximum emission region size can be associated with a light crossing time of a few months or less. For typical Galactic distances of a few kiloparsecs, this implies angular extents on the sky $\lesssim 1'$. An acceleration mechanism that can generate very high energy particles in a short period must also be present. One known source class fitting these criteria are pulsar wind nebulae (PWN).

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2. The Candidate PWN

There are around 10 X-ray and radio PWN discovered so far around known pulsars (eg. Gaensler et al. 2000; Chevalier 2000), with a similar number of good candidates without a detected pulsar. Remarkably, around half of these sources are associated with γ -ray pulsars or coincident with unidentified EGRET sources. Only the oldest γ -ray pulsars, such as Geminga and PSR B1055-52, show no sign of having PWN. Therefore, it might be expected that other γ -ray pulsars might reveal themselves through PWN.

None of the known γ -ray pulsars are significantly variable as measured above 100 MeV using the τ statistic (Tompkins 1999). However, in the 70-150 MeV range, the unpulsed flux of the Crab was observed to vary (de Jager et al. 1996), probably a synchrotron emission component from the inner nebula, perhaps associated with the variable optical and X-ray wisps (Scargle 1967, and Hester etc. these proceedings), with a variable exponential cutoff at around 25 MeV. Chandra has also revealed spatial variability of the X-ray emission on a timescale of a few months in the Vela and G11.2-0.3 PWN (Kargaltsev et al., Roberts et al. these proceedings). Therefore, high-energy variability on the appropriate time and spatial scales appears to be a common feature of PWN, and the question becomes that of what energy the γ -ray spectral cut-off is and why.

In a recent flux-limited survey of the brightest sources of emission above 1 GeV, Roberts, Romani, and Kawai (2001) noted that the four unidentified sources showing the most variability above 100 MeV (out of 20 galactic sources in their sample measured for variability by Tompkins 1999) contained the brightest new sources of extended hard X-ray emission found in the survey and the previously known PWN around PSR B1853+01 in the supernova remnant W44. The supernova remnant W44 is one of a class of remnants which is centrally filled with thermal X-rays, implying a relatively dense interior. The radio PWN around PSR B1853+01 is a plume reaching only to the north of the pulsar, suggesting a ram-pressure confined wind (Slane et al. 1997). The hard X-ray emission is also offset to the north.

A second variable source, GeV J1417–6100, coincides with two extended hard X-ray sources. Radio studies of these sources show they are within a complex of radio emission called the Kookaburra. One is coincident with the Rabbit radio nebula, whose morphology, relatively flat (spectral index $\alpha = -0.25$) spectrum, and high polarization all mark it as a probable PWN (Roberts et al. 1999). The polarized emission has a trail-like morphology, the head of which is coincident with an X-ray peak seen in the ASCA SIS. The second X-ray source is at the southern edge of the upper wing of the Kookaburra, whose high ratio of infrared to radio emission, excess polarization, and non-thermal spectrum mark it as a likely supernova remnant, although whether it is partially or fully a result of a pulsar wind is difficult to determine with the existing data (Roberts, Romani & Johnston 2001). This second X-ray source has recently been found to contain the high \dot{E} pulsar PSR J1420–6048 (D'Amico et al. 2001), and is therefore an X-ray and probably radio PWN.

A third variable source, GeV J1825–1310, contains a newly discovered extended hard X-ray source with no obvious bright radio counterpart. We have

imaged the region with the VLA at 20 cm, and find much excess radio emission in the region, although precise measurements are difficult due to several nearby bright HII regions. The MSX satellite 8.3 $\mu{\rm m}$ image shows the infrared emission in the region to be partially correlated with the radio emission, but not clearly so in the vicinity of the X-ray nebula. Although the X-ray nebula in GeV J1825–1310 does not show all of the observational signatures of a PWN, other interpretations of its nature seem even less well justified.

The fourth variable source, GeV J1809–2327, contains an extended X-ray complex partially coincident with stars in an OB association. It is contained within the Lynds 227 dark nebula, and appears to be interacting with a cloud imaged in CO (Oka et al. 1999). A recent short Chandra ACIS exposure resolves much of the emission into point sources coincident with stars in the OB association and embedded in a thermal radio cloud, but it also reveals a point source embedded in diffuse emission somewhat to the North of the stellar cluster with no optical counterpart but with a short bright trail of emission leading from it (Braje et al. 2001). This latter source is at the tip of a funnel-shaped radio nebula whose high polarization and spectral index strongly suggest it is a PWN (Roberts, Gaensler & Romani 2002). With this GeV source also coinciding with a likely X-ray and radio PWN, we are led to the hypothesis that some of the variable EGRET sources are PWN.

3. Emission Mechanisms

de Jager et al. (1996) argue that a general upper limit to the cut-off energy of synchrotron emission from shock accelerated electrons can be derived by balancing the acceleration timescale (equal to the electron gyroperiod across a strong shock) with the synchrotron loss timescale. The resulting cutoff energy is:

$$E_{max} = 25(D\alpha/\sin\theta) \,\text{MeV}$$

Here D is the Doppler boost factor, θ is the electron pitch angle, and $\alpha < 1$. If this is the true limit on emission from shock generated electrons, then to reach energies of several hundred MeV, either the Doppler factor is large or the pitch angle is small. If the one-sided nature of some of the nebulae is due to large Doppler boost factors then perhaps the emission can reach the required energies.

On the other hand, Gallant & Arons (1994) suggest that ions play an important role in the generation of wisps, and the spatial and time scales are related to the ion cyclotron radius, with emission coming from electrons coupled to the ions. For ion-pumped shock acceleration, the maximum emission energy is proportional to the ambient magnetic field $E_{max} \propto B\gamma^2 \propto B\dot{E}$. Assuming that the variable γ -ray emission is synchrotron radiation, we can estimate a lower limit on the magnetic field in the emitting region. It is reasonable to assume that the synchrotron cooling timescale must be shorter than the variability time scale $(t_{cool} \leq t_{var})$. Therefore, a lower limit on the magnetic field is given by:

$$B \ge 10^{-4} \left(\frac{10 \text{MeV}}{E}\right)^{1/3} t_{var}^{-2/3}$$

For E = 100 MeV and $t_{var} = 0.2-2$ yr., the magnetic field is $B > 0.3-1.4 \times 10^{-4}$ G, similar to the inferred magnetic fields in the Crab and Vela PWN.

Observationally, the γ -ray emitting candidates seem to be in dense environments (near or in molecular clouds or overdense SNR), and may all be bow-shock nebulae. The magnetic field could then be compressed and amplified if the pulsars have high velocities:

$$E_{max} \propto B\dot{E} \propto \sqrt{\rho}v\dot{E}$$

although to increase the cutoff to many times that of the Crab from pulsars whose spin-down energy is one or two orders of magnitude smaller may be quite difficult.

Alternatively, perhaps other emission mechanisms are important. Bremm-strahlung radiation from the pulsar wind striking an ambient molecular cloud (Oka et al. 1999) has been suggested in the case of GeV J1809-2327, although it may be too inefficient if we require the emission region to be as small as implied by the variability time scale. Inverse Compton scattering is also an attractive alternative. However, the cooling timescale of the electrons involved for IC emission in the EGRET waveband is much longer than the variability timescale, implying the density of low-energy photons that are upscattered at the emission site would need to be modulated in some manner.

Ultimately, GLAST and AGILE will make unambiguous source identifications, and will be able to verify the variability seen by EGRET, allowing correlations to be made with the low-energy spatial variability. However, if a significant fraction of the emission is truly variable, the pulse fraction of these sources is small, and the detection of any pulsed γ -ray emission through blind searches is unlikely to succeed. Therefore, X-ray and radio pulse detections will still be necessary to detect γ -ray pulsations and determine the energetics of these systems.

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